

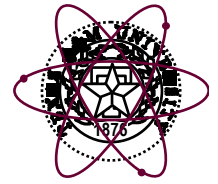
Quadrature Class

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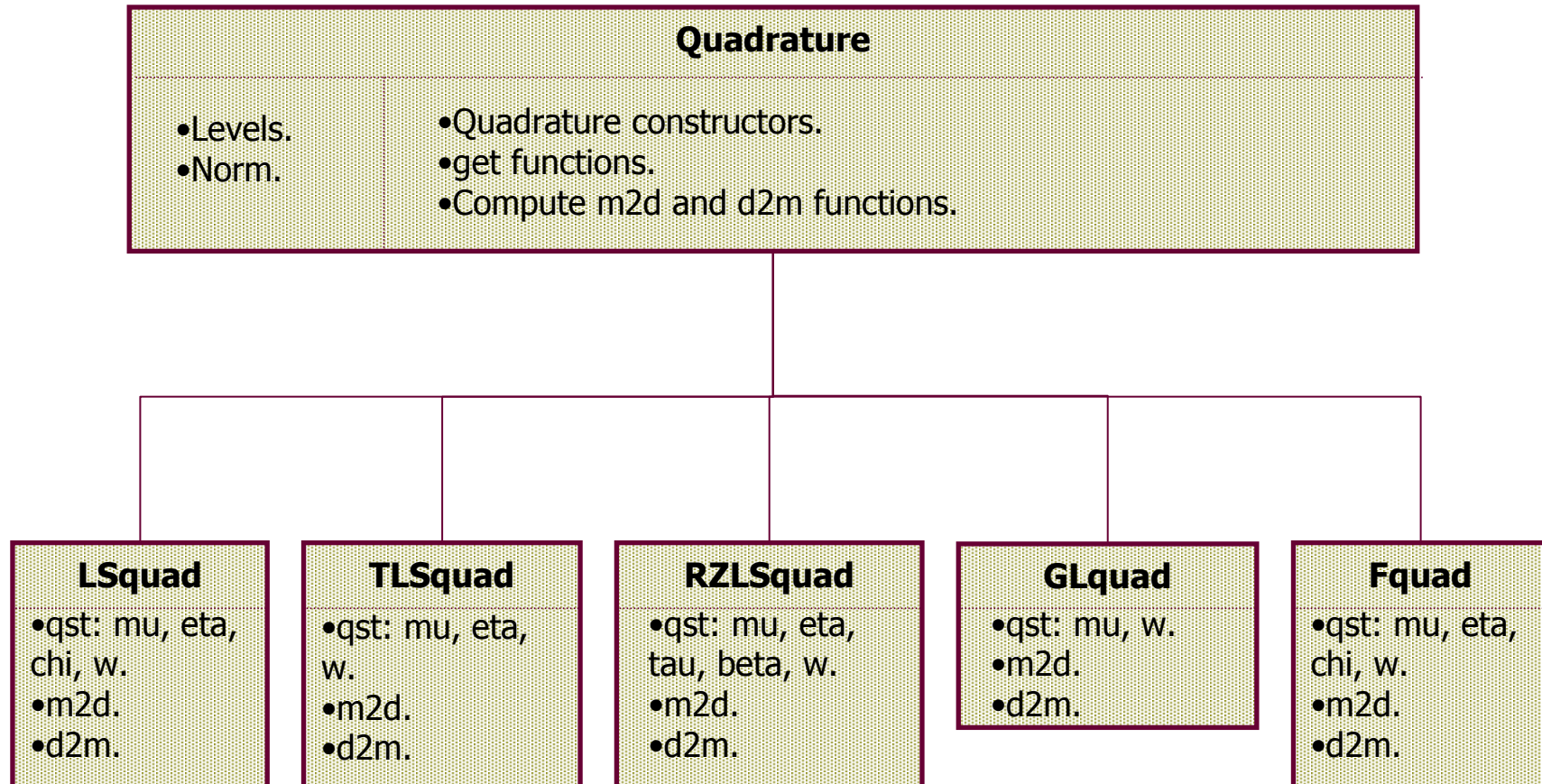
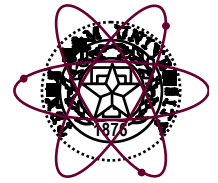
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We have group angular discretization quadratures into a single class.

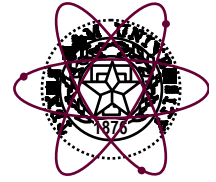


- Primarily, we store quadrature points, weights and norms.
 - ⇒ *Struct contains each quadrature point and weight.*
 - ⇒ *Vector of struct groups information for all directions.*
- We store the expansion coefficient for the construction of flux moments based on angular flux information (discrete to moment) and vice versa (moment to discrete).
 - ⇒ *Two vector of vectors contain all coefficients for each direction: d_{2m} and m_{2d} .*
 - ⇒ *Weights are build using only the real components of the spherical harmonic functions.*
- This m_{2d} and d_{2m} approach is beneficial in two ways:
 - ⇒ *m_{2d} and d_{2m} are only constructed once.*
 - ⇒ *We can expand easily to Garlekin Quadratures.*

We organize the quadrature class hierarchy based on the problem's geometry.

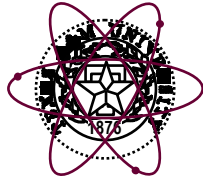


We construct the quadrature by storing magnitudes and exploiting symmetries.



- We store the magnitudes used in each quadrature point for quadratures with 2 to 24 levels.
- We assign the quadrature points and weights to the 1st octant.
- We expand to each octant using the assumed symmetries.
- There are three special cases:
 - ⇒ *RZ Level symmetric quadrature, assigns the quadrature points and calculates the angle differencing coefficients in order.*
 - ⇒ *Gauss-Legendre quadrature calculates the quadrature points and weights.*
 - ⇒ *Filtering quadrature always uses 6 points lying on the cartesian axis.*

We construct the m2d and d2m vectors using spherical harmonic expansions.



- Based on Kevin Clarno's work on anisotropic scattering.

$$\int_{4\pi} d\Omega \Sigma_s(\underline{r}, E' \rightarrow E, \underline{\Omega}' \bullet \underline{\Omega}) \psi(\underline{r}, E', \underline{\Omega}') = \left\{ \begin{array}{l} \int_{4\pi} d\Omega' \left[\sum_{k=0}^K \frac{2k+1}{norm} \Sigma_{s,k}(\underline{r}, E' \rightarrow E) P_k(\mu_0) \right] \\ \left[\sum_{l=0}^K \frac{2l+1}{norm} \sum_{m=-l}^l \phi_{lm}(\underline{r}, E') Y_{lm}(\underline{\Omega}') \right] \end{array} \right\}$$

- After applying the addition theorem :

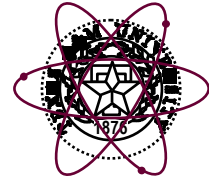
$$S_{\text{scatt}}(\underline{r}, E, \underline{\Omega}) = \sum_{k=0}^K \frac{2k+1}{norm} \Sigma_{s,k}(\underline{r}, E' \rightarrow E) \sum_{n=-k}^k \phi_{kn}(\underline{r}, E') Y_{kn}(\underline{\Omega})$$

$$S_{\text{scatt}}(\underline{r}, E, \underline{\Omega}) = \sum_{k=0}^K \frac{2k+1}{norm} \Sigma_{s,k}(\underline{r}, E' \rightarrow E) \left\{ \begin{array}{l} Y_{k0}(\underline{\Omega}) \phi_{k0}(\underline{r}, E') + \\ \sum_{n=1}^k \left[\phi_{kn}(\underline{r}, E') Y_{kn}(\underline{\Omega}) + \right. \\ \left. \phi_{k-n}(\underline{r}, E') Y_{k-n}(\underline{\Omega}) \right] \end{array} \right\}$$

- Based on the spherical harmonic definitions:

$$\left. \begin{array}{l} Y_{k,-n}(\underline{\Omega}) = (-1)^n Y_{kn}^*(\underline{\Omega}) \\ \phi_{k,-n}(\underline{r}, E') = (-1)^n \phi_{kn}^*(\underline{r}, E') \end{array} \right\} \Rightarrow S_{\text{scatt}}(\underline{r}, E, \underline{\Omega}) = \sum_{k=0}^K \frac{2k+1}{norm} \Sigma_{s,k}(\underline{r}, E' \rightarrow E) \left\{ \begin{array}{l} Y_{k0}(\underline{\Omega}) \phi_{k0}(\underline{r}, E') + \\ \sum_{n=1}^k \left[\phi_{kn}(\underline{r}, E') Y_{kn}(\underline{\Omega}) + \right. \\ \left. \phi_{kn}^*(\underline{r}, E') Y_{kn}^*(\underline{\Omega}) \right] \end{array} \right\}$$

The spherical harmonics functions can be expressed solely with real numbers.



$$f + f^* = 2 \operatorname{Re}[f] \Rightarrow S_{\text{scatt}}(\underline{r}, E, \underline{\Omega}) = \sum_{k=0}^K \frac{2k+1}{\text{norm}} \Sigma_{s,k}(\underline{r}, E' \rightarrow E) \left\{ \begin{array}{l} Y_{k0}^{\operatorname{Re}}(\underline{\Omega}) \phi_{k0}^{\operatorname{Re}}(\underline{r}, E') \\ + 2 \sum_{n=1}^k \operatorname{Re}[\phi_{kn}(\underline{r}, E') Y_{kn}(\underline{\Omega})] \end{array} \right\}$$

- **Rearranging:**

$$S_{\text{scatt}}(\underline{r}, E, \underline{\Omega}) = \sum_{k=0}^K \frac{2k+1}{\text{norm}} \Sigma_{s,k}(\underline{r}, E' \rightarrow E) \left\{ \begin{array}{l} Y_{k0}^{\operatorname{Re}}(\underline{\Omega}) \phi_{k0}^{\operatorname{Re}}(\underline{r}, E') \\ + 2 \sum_{n=1}^k \left[\begin{array}{l} \operatorname{Re}[\phi_{kn}(\underline{r}, E')] \operatorname{Re}[Y_{kn}(\underline{\Omega})] \\ - \operatorname{Im}[\phi_{kn}(\underline{r}, E')] \operatorname{Im}[Y_{kn}(\underline{\Omega})] \end{array} \right] \end{array} \right\}$$

- **We define the variables:**

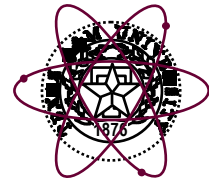
$$R_{kn} = \begin{cases} \sqrt{2 - \delta_{n0}} \operatorname{Re}[Y_{kn}] \forall n \geq 0 \\ \sqrt{2} \operatorname{Im}[Y_{k,n}] \forall n < 0 \end{cases}, \tilde{\phi}_{kn}(\underline{r}, E') = \int_{4\pi} d\Omega' [R_{kn}(\Omega') \psi(\underline{r}, E', \Omega')]$$

- **Applying these definition to the discrete case, d2m and m2d matrices become:**

$$S_{\text{scatt}}(\underline{r}, E, \underline{\Omega}) = \sum_{k=0}^K \sum_{n=-k}^k \left[\frac{2k+1}{\text{norm}} R_{kn}(\Omega_m) \right] \Sigma_{s,k}(\underline{r}, E_g' \rightarrow E_g) \tilde{\phi}_{kn}(\underline{r}, E_g)$$

$$\tilde{\phi}_{kn}(\underline{r}, E_g) = \sum_{m=1}^M w_m R_{kn}(\Omega_m) \psi(\underline{r}, E_g, \Omega_m)$$

More work needs to be done ...



- **The test passes isotropic scattering:**
 - ⇒ *Source iterations test.*
 - ⇒ *GMRES.*
 - ⇒ *TSA , Stretch and Filter.*
- **We are searching for tests for anisotropic scattering in xyz and xy geometries.**
- **We need to change our inputs:**
 - ⇒ *Currently only scattering order is specified at the base of the XML tree input.*
 - ⇒ *User will specify the scattering order per isotope.*
 - The maximum scattering order over all isotopes will be the scattering order.
 - ⇒ *User will specify type of scattering expansion.*